

# Effects of Nb<sub>2</sub>O<sub>5</sub> doping on the microstructure and the dielectric temperature characteristics of barium titanate ceramics

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**Abstract** In this work, the effects of Nb<sub>2</sub>O<sub>5</sub> addition on the dielectric properties and phase formation of BaTiO<sub>3</sub> were investigated. A core–shell structure was formed for Nb-doped BaTiO<sub>3</sub> resulted from a low diffusivity of Nb<sup>5+</sup> ions into BaTiO<sub>3</sub> when grain growth was inhibited. In the case of 0.3–4.8 mol% Nb<sub>2</sub>O<sub>5</sub> additions, two dielectric constant peaks were observed. The Curie dielectric peak was determined by the ferroelectric–paraelectric transition of grain core, whereas the secondary broad peak at lower temperature was due to strong chemical inhomogeneity in Nb-doped BaTiO<sub>3</sub> ceramics. The dielectric constant peak at Curie temperature was markedly depressed with the addition of Nb<sub>2</sub>O<sub>5</sub>. On the other hand, the secondary dielectric constant peak was enhanced when sintered above 1280 °C for higher Nb<sub>2</sub>O<sub>5</sub> concentrations (≥1.2 mol%). The Curie temperature was shifted to higher temperatures, whereas the transition temperature corresponding to the secondary peak moved to lower temperatures as increasing the amount of Nb<sub>2</sub>O<sub>5</sub> more than 1.2 mol%. The decrease of this lower transition temperature was assumed to be closely related with the secondary phase formation when Nb concentration greater than 1.2 mol%. From XRD analyses, a large amount of secondary phases was observed when Nb<sub>2</sub>O<sub>5</sub> amount exceeded 1.2 mol%. The coefficients of thermal expansion of Nb-doped BaTiO<sub>3</sub> were increased with increasing Nb<sub>2</sub>O<sub>5</sub> contents, resulting in large internal stress between cores and shells. Therefore, the shift of Curie temperature to higher temperatures was attributed to

internal stress resulting from the formation of a core–shell structure and a large amount of secondary phase grains.

## Introduction

Barium titanate (BaTiO<sub>3</sub>, BT) is well known as one of the most widely studied and used perovskite ferroelectrics. It has a ferroelectric tetragonal structure at room temperature and changes to a paraelectric cubic structure on heating. The ferroelectric phase-transition temperature, called as Curie temperature or  $T_C$ , is a key factor for preparation of BaTiO<sub>3</sub>-based temperature-stable capacitors because dielectric constant will decrease sharply above  $T_C$  according to the Curie–Weiss law  $\varepsilon = C/(T - T_C)$  [1]. Hence, increasing the Curie temperature of BaTiO<sub>3</sub> is favorable in improving the temperature stability of dielectric constant so as to meet the requirement of high-temperature capacitors.

The Curie temperature is affected by such many factors as modification of chemical composition [2–5], particle and grain size [6, 7], tetragonality [7, 8], oxygen vacancies concentration [9, 10], stress [4, 5, 11, 12], and so on. The doping effects on the  $T_C$  of BaTiO<sub>3</sub> have been extensively investigated. As is well known, most impurity additions lower the Curie temperature of BaTiO<sub>3</sub>. The Nb-doping effect on the Curie temperature of BaTiO<sub>3</sub> has been reported by many researchers. Kahn reported that the Curie temperature of BaTiO<sub>3</sub> was shifted to lower temperatures by Nb addition when significant grain growth happened. Contrarily, the Curie temperature of BaTiO<sub>3</sub> moved to higher temperatures for small-grained BaTiO<sub>3</sub> ceramics when grain growth was inhibited [13]. Cui et al. [14] reported that the Curie temperature of BaTiO<sub>3</sub> was

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decreased by Nb addition. In view of these inconsistent results, the present study focused on the investigation of the effect of Nb addition on the dielectric properties of BaTiO<sub>3</sub>. The results proved that Nb shifted the Curie temperature to higher temperatures. This paper discusses the influence of Nb-doping on the Curie temperature of BaTiO<sub>3</sub> and possible mechanism.

## Experimental procedure

Samples were prepared from commercial BaTiO<sub>3</sub> powder (GuoTeng Ceramic Inc., average grain size 0.4 μm, Ba/Ti = 1) doped with 0.3, 0.6, 1.2, 2.4, 3.6, or 4.8 mol% Nb<sub>2</sub>O<sub>5</sub> (99% purity) by the conventional ceramic processing technique. Then, 1 mol% BaSiO<sub>3</sub> was used as sintering aid. Raw materials were ball-milled in deionized water for 8 h. The prepared ceramic powders were pressed in disk form (10 mm in diameter and 1 mm thick) by mixing ceramic powders with 3 wt% PVA binder. After debinding, the disks were finally fired at 1240–1320 °C for 2 h in air. Fired-on-silver was used as electrodes for the measurement of electric properties. Dielectric measurement of samples was performed by using a LCR meter (YY2812) auto-controlled by computer at 1 kHz and 1.0 V rms in the temperature range from –55 to 200 °C. Microstructures of the ceramics were observed using an S-530 scanning electron microscope (SEM). The X-ray diffraction (XRD) analysis was carried out using a Philips X'Pert diffractometer over the range of  $2\theta = 20\text{--}60^\circ$  with CuK $\alpha$  radiation.

The recently developed software Material Analysis Used Diffraction (MAUD), which is based on a Rietveld method combined with a Fourier analysis, has been applied to analyze the XRD of alloys and ferroelectrics [4, 15–17]. In our work, the structural refinement was carried out by using MAUD. The coefficients of thermal expansion for pure BaTiO<sub>3</sub> and Nb-doped BaTiO<sub>3</sub> ceramics were measured by thermal expansion meter with a heating rate of 5 °C/min. The coefficient of thermal expansion was calculated by Eq. 1:

$$\bar{\alpha} = \frac{1}{L_0} \cdot \frac{L_2 - L_1}{T_2 - T_1} + \alpha_0 \quad (1)$$

where  $L_0$  is the length of sample at room temperature,  $L_1$ ,  $L_2$  is the length of sample at  $T_1$  and  $T_2$ , respectively, and  $\alpha_0$  is the coefficient of thermal expansion of quartz ( $5.7 \times 10^{-7}/^\circ\text{C}$ ).

## Results and discussion

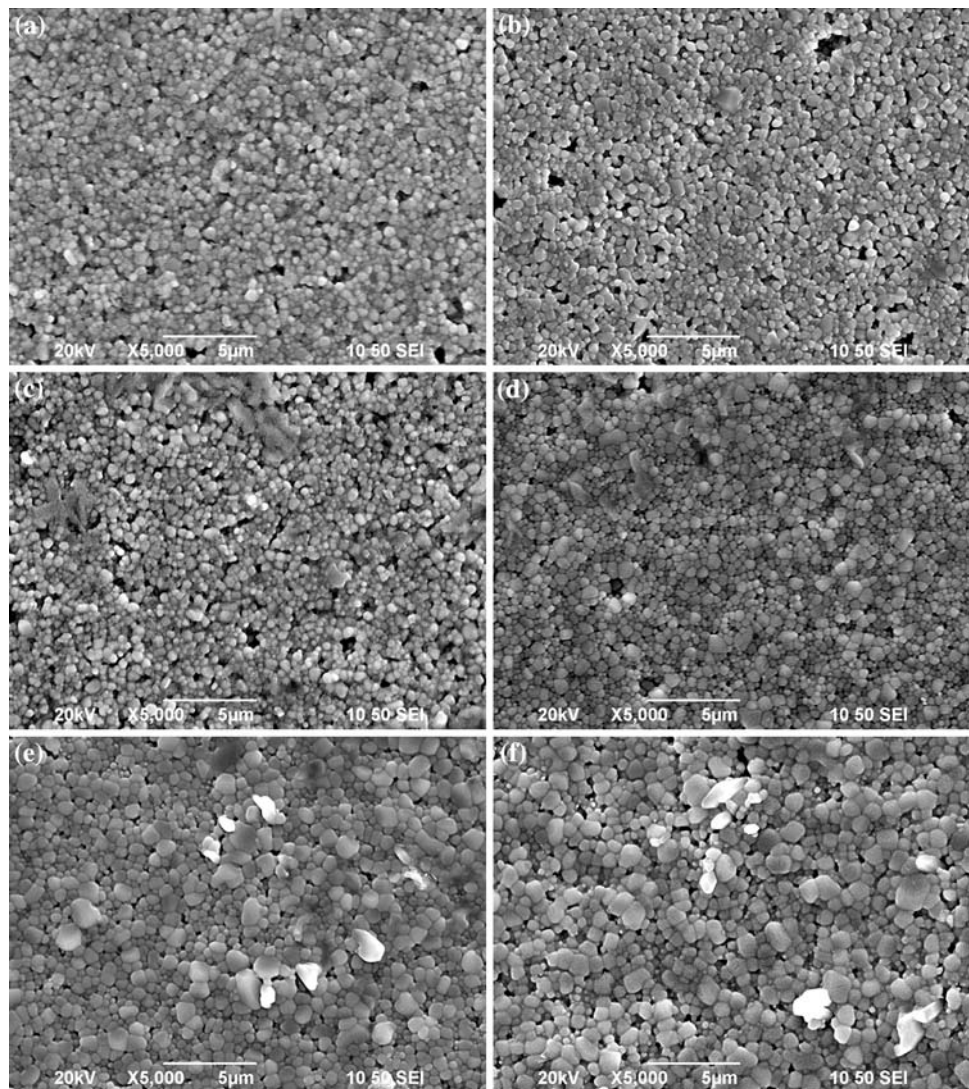
SEM micrographs for samples sintered at 1280 °C with various Nb<sub>2</sub>O<sub>5</sub> contents were shown in Fig. 1. When

0.3–2.4 mol% Nb<sub>2</sub>O<sub>5</sub> was added, the microstructure was homogeneous fine grained and rather porous (Fig. 1a–d). As Nb<sub>2</sub>O<sub>5</sub> concentration increased further ( $\geq 3.6$  mol%), significant grain growth occurred and the proportion of pores decreased (Fig. 1e, f). As Nb<sub>2</sub>O<sub>5</sub> content increased to 1.2 mol%, secondary phase began to occur which was proved by XRD analysis. However, it was hard to observe the secondary phase grains by SEM due to the formation of small amount of secondary phase. Figure 2 shows the grain morphology of samples sintered at 1320 °C doped with 0.3, 1.2 and 2.4 mol% Nb<sub>2</sub>O<sub>5</sub>, respectively. It was observed that sintering temperature had little effect on the size of matrix grains for Nb-doped BaTiO<sub>3</sub> ceramics. As 0.3–1.2 mol% Nb<sub>2</sub>O<sub>5</sub> was doped, no secondary phase appeared (Fig. 2a, b). With 2.4 mol% Nb<sub>2</sub>O<sub>5</sub> addition, two types of grains were observed, that is, the matrix composed of fine grains and large secondary phase grains (Fig. 2c).

The phase formation of samples sintered at 1280 °C with different Nb<sub>2</sub>O<sub>5</sub> contents (0.6–4.8 mol%), as examined by XRD is shown in Fig. 3. It was obvious that the perovskite BaTiO<sub>3</sub> phase was formed in all samples. With an addition of 0.6 mol% Nb<sub>2</sub>O<sub>5</sub>, no additional peak representing a secondary phase was observed. However, secondary phases were detected when Nb<sub>2</sub>O<sub>5</sub> amount exceeded 1.2 mol%. The composition of the secondary phases was identified as Ba<sub>6</sub>Ti<sub>14</sub>Nb<sub>2</sub>O<sub>39</sub>, Ba<sub>3</sub>Ti<sub>5</sub>Nb<sub>3,2</sub>O<sub>21</sub>, and Ba<sub>6</sub>Ti<sub>17</sub>O<sub>40</sub>. Therefore, it was assumed that the solubility limit of Nb<sub>2</sub>O<sub>5</sub> in BaTiO<sub>3</sub> at 1280 °C in air was less than 1.2 mol%. The secondary phase formation was due to the substitution of Nb<sup>5+</sup> to Ti<sup>4+</sup> site and segregation of Ti<sup>4+</sup> out of BaTiO<sub>3</sub> grains [18].

Figure 4a shows the temperature dependence of dielectric constant for BaTiO<sub>3</sub> ceramics sintered at 1240 °C with various Nb<sub>2</sub>O<sub>5</sub> contents. With 0.3 mol% Nb<sub>2</sub>O<sub>5</sub> addition, two dielectric constant peaks were observed. A core–shell structure was developed for small grained BaTiO<sub>3</sub> ceramics during sintering because of low diffusivity of Nb<sup>5+</sup> ions into BaTiO<sub>3</sub> lattice [19]. The core was composed of pure barium titanate, while the shell was doped with additive Nb. The maximum of dielectric constant at 133 °C was determined by the phase transition from ferroelectric to paraelectric of unreacted pure BaTiO<sub>3</sub> grain core so that this temperature was called Curie temperature or  $T_C$ . On the other hand, the dielectric constant peak at a lower temperature of 45 °C was broad and diffuse due to strong chemical inhomogeneity in Nb-doped BaTiO<sub>3</sub> ceramics. The transition temperature corresponding to the secondary broad peak was denoted as  $T_1$ . With an increase of Nb<sub>2</sub>O<sub>5</sub> content, the dielectric constant peak at  $T_C$  was obviously depressed, especially in the cases of  $\geq 1.2$  mol% Nb<sub>2</sub>O<sub>5</sub> additions. It is also shown in Fig. 4a that  $T_1$  shifts to lower temperatures as Nb amount exceeded 1.2 mol%. With  $>3.6$  mol% Nb<sub>2</sub>O<sub>5</sub> addition, the secondary dielectric constant peak even could not be

**Fig. 1** SEM micrographs of  $\text{Nb}_2\text{O}_5$ -doped  $\text{BaTiO}_3$  ceramics sintered at 1280 °C.  $\text{Nb}_2\text{O}_5$  contents of these samples were **a** 0.3, **b** 0.6, **c** 1.2, **d** 2.4, **e** 3.6, and **f** 4.8 mol%

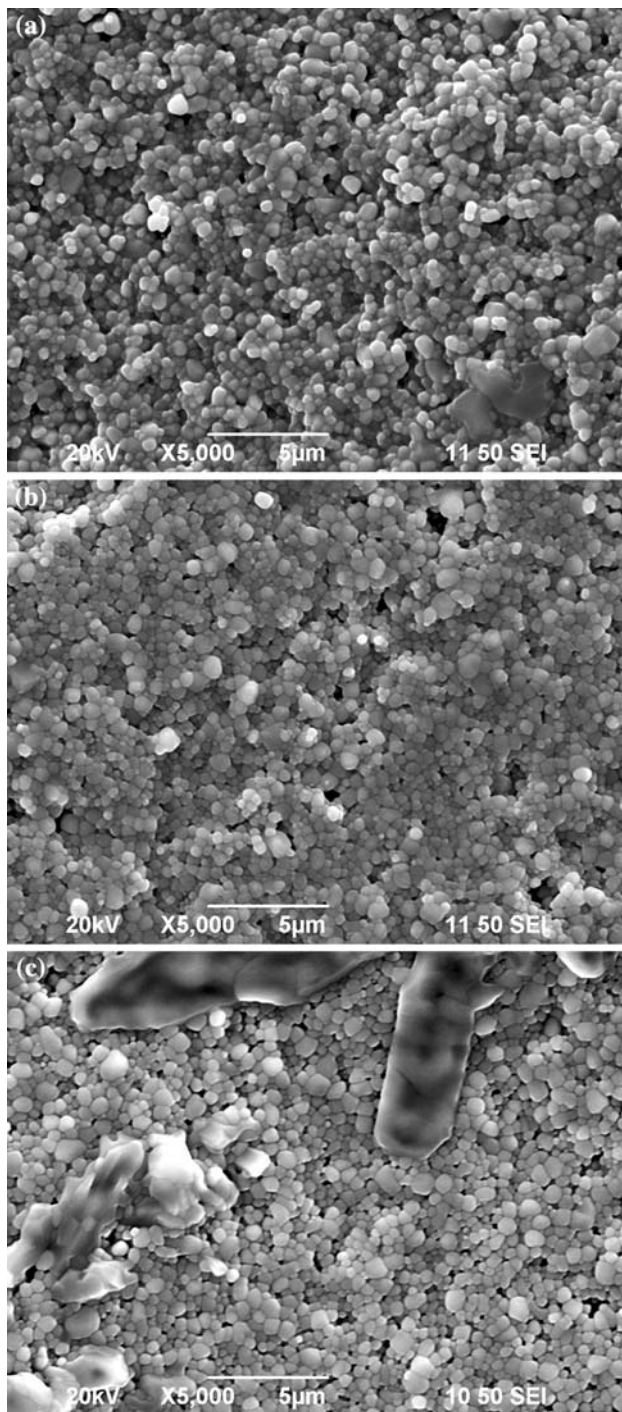


observed in the temperature range of  $-55$  to  $200$  °C. This indicated that Nb made  $T_1$  lower than  $-55$  °C for higher  $\text{Nb}_2\text{O}_5$  concentrations ( $>3.6$  mol%). Figure 4b and c shows the temperature dependence of dielectric constant for Nb-doped  $\text{BaTiO}_3$  ceramics sintered at 1280 and 1320 °C with various  $\text{Nb}_2\text{O}_5$  contents, respectively. It is obvious that the curves in Fig. 4b and c are similar to those in Fig. 4a. The impact of sintering temperature on the temperature dependence of dielectric constant is slight for the sample doped with 0.3 mol%  $\text{Nb}_2\text{O}_5$ , which suggested that the grain shell is thin when a low concentration  $\text{Nb}_2\text{O}_5$  was added. However, as  $\text{Nb}_2\text{O}_5$  concentration exceeded 1.2 mol%, the dielectric constant peak at  $T_C$  was greatly depressed, whereas the dielectric constant peak at lower temperature was enhanced with increasing sintering temperature. Especially, for the samples doped with more than 3.6 mol%  $\text{Nb}_2\text{O}_5$ , the Curie peak disappeared by sintering at temperature above 1280 °C, indicating the collapse of core-shell structure and homogeneous distribution of Nb in  $\text{BaTiO}_3$ .

High sintering temperature promoted solid-state diffusion and caused the core-shell structure to disappear.

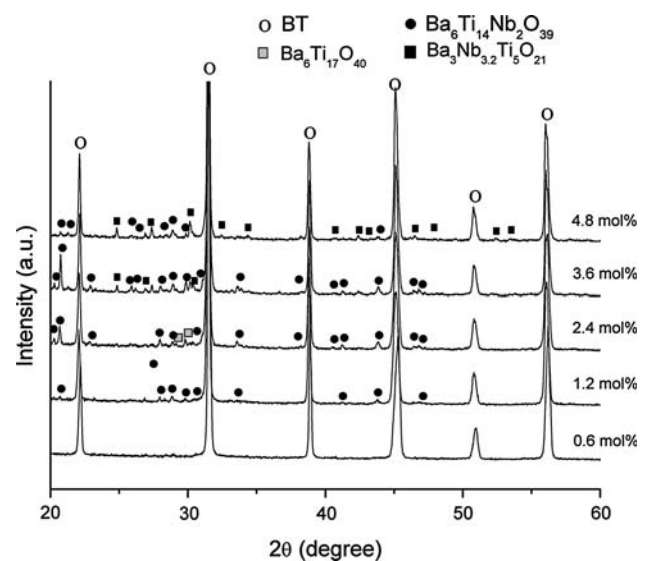
Figure 5 shows the dependences of  $T_C$  and  $T_1$  on  $\text{Nb}_2\text{O}_5$  contents for Nb-doped  $\text{BaTiO}_3$  ceramics sintered at 1240, 1280, or 1320 °C, respectively. With increasing the amount of  $\text{Nb}_2\text{O}_5$ ,  $T_C$  increased, while  $T_1$  was independent of Nb content up to 1.2 mol% and then decreased rapidly with further addition. It was observed that  $T_1$  was almost unchanged for lower  $\text{Nb}_2\text{O}_5$  concentrations ( $\leq 1.2$  mol%). Moreover, sintering temperature has no effect on  $T_1$  at  $\text{Nb}_2\text{O}_5$  content less than 1.2 mol%. It has been known that high sintering temperature promoted solid-state diffusion of Nb and caused the volume fraction of grain shell to increase. Therefore, it was assumed that  $T_1$  had little to do with the volume fraction of grain shell. For higher  $\text{Nb}_2\text{O}_5$  concentrations ( $>1.2$  mol%),  $T_1$  was related with Nb content, that is,  $T_1$  was almost decreased linearly as  $\text{Nb}_2\text{O}_5$  increased. Moreover,  $T_1$  shifted to lower temperatures by increasing sintering temperature from 1240 to 1280 °C. It





**Fig. 2** SEM micrographs of  $\text{Nb}_2\text{O}_5$ -doped  $\text{BaTiO}_3$  ceramics sintered at  $1320^\circ\text{C}$ .  $\text{Nb}_2\text{O}_5$  contents of these samples were **a** 0.3, **b** 1.2, and **c** 2.4 mol%

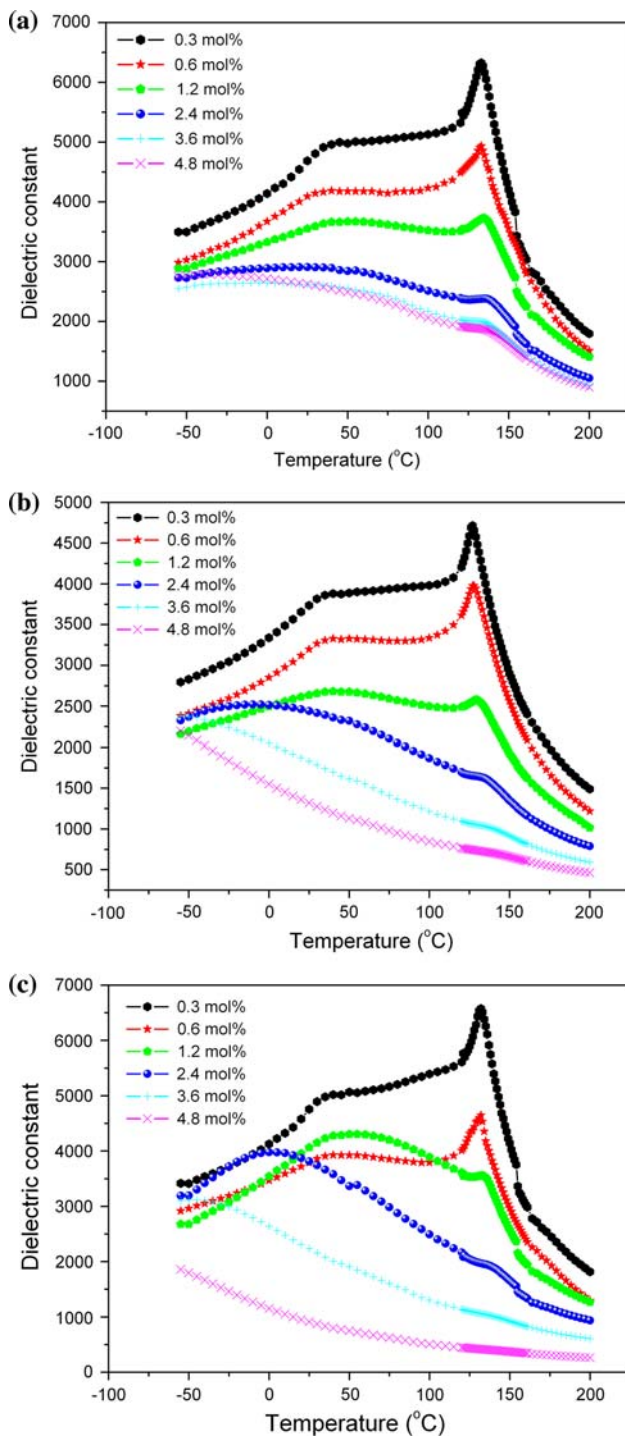
was analyzed by XRD that the solubility limit of  $\text{Nb}_2\text{O}_5$  in  $\text{BaTiO}_3$  was less than 1.2 mol%. At  $\text{Nb}_2\text{O}_5$  content above 1.2 mol%, secondary phase appeared and the amount of the secondary phase increased with an increase in  $\text{Nb}_2\text{O}_5$  content and sintering temperature. Therefore, the decrease



**Fig. 3** X-ray diffraction patterns of samples with various  $\text{Nb}_2\text{O}_5$  contents

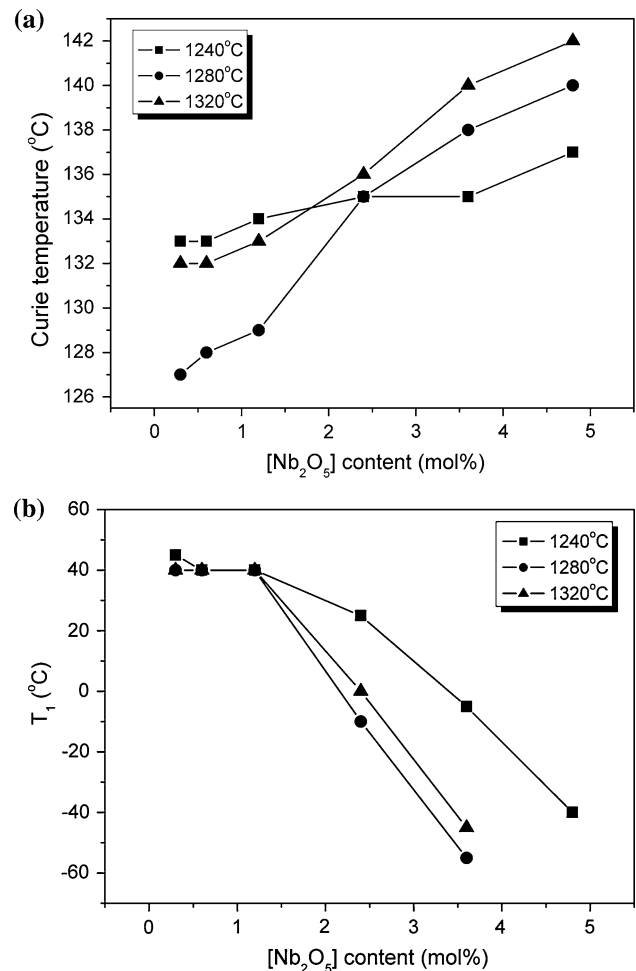
of  $T_1$  was assumed to be correlated with the secondary phase formation when  $\text{Nb}_2\text{O}_5$  concentration is greater than 1.2 mol%.

In contrast to  $T_1$ ,  $T_C$  shifted to higher temperatures with  $\text{Nb}_2\text{O}_5$  addition in the present study. According to ionic size,  $\text{Nb}^{5+}$  ions preferentially substituted to  $\text{Ti}^{4+}$  and extra positive charge could be compensated by  $V''_{\text{Ba}}$  and/or  $V''''_{\text{Ti}}$ . It was well accepted that the coupling of  $\text{TiO}_6$  oxygen octahedral played a critical role in the stability of ferroelectric phase in  $\text{BaTiO}_3$  perovskite [20]. The degree of the coupling between neighboring  $\text{TiO}_6$  octahedra or the stability of ferroelectric phase would be significantly weakened by introducing defects. Therefore, Nb addition would shift the Curie temperature of  $\text{BaTiO}_3$  to lower temperatures. Kahn [13] reported that incorporation of single Nb decreased the Curie temperature in large-grained  $\text{BaTiO}_3$  and small-grained  $\text{BaTiO}_3$  grown from ultrafine powder, in agreement also with results by Hennings and Rosenstein [3] and Cui et al. [14]. During sintering,  $\text{Nb}^{5+}$  ions diffused into  $\text{BaTiO}_3$  lattice and a core-shell structure would be formed due to a low diffusivity of  $\text{Nb}^{5+}$ . However, the diffusion of  $\text{Nb}^{5+}$  ions was enhanced and the core-shell structure would be destroyed as significant grain growth happened, which led to a homogeneous distribution of Nb in  $\text{BaTiO}_3$  [13, 21]. We attributed the decrease of  $T_C$  in their work to a homogeneous distribution of Nb in  $\text{BaTiO}_3$ . In our study, a core-shell structure was formed, which was proved by the double dielectric constant peak in Fig. 4. The shift of Curie temperature to higher temperatures resulted from the internal stress existing in core-shell structure grains. As discussed earlier,  $T_C$  was correlated with the ferroelectric-paraelectric phase transition of grain core. Jung et al. [5], Song et al. [12], and Hwang et al. [22] have suggested that



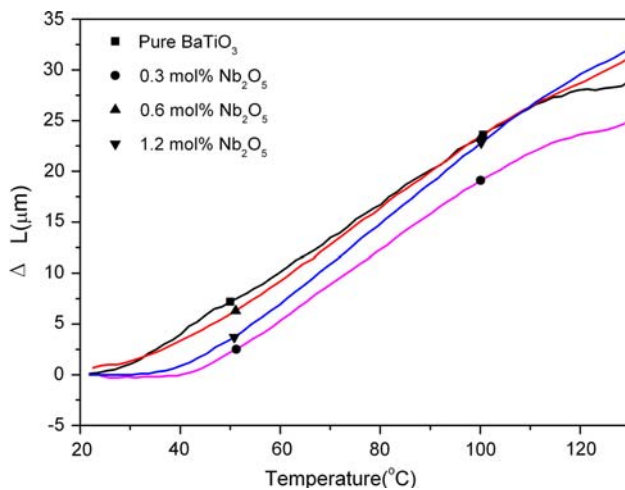
**Fig. 4** Temperature dependence of dielectric constant for BaTiO<sub>3</sub> ceramics with various Nb<sub>2</sub>O<sub>5</sub> contents: **a** 1240 °C, **b** 1280 °C, and **c** 1320 °C

the addition of smaller ionic radius rare earth elements moved  $T_C$  to higher temperatures in samples composed of a core-shell structure in the BaTiO<sub>3</sub>-MgO-Re<sub>2</sub>O<sub>3</sub> (Re = Ho, Er, Yb, Lu, Y) system. The dielectric constant peak at  $T_C$  was substantially determined by the



**Fig. 5** Curie temperature  $T_C$  and  $T_1$  as functions of Nb<sub>2</sub>O<sub>5</sub> contents for BaTiO<sub>3</sub> ceramics sintered at 1240, 1280, and 1320 °C, respectively

cubic-tetragonal phase transition of unreacted pure BaTiO<sub>3</sub> grain core. They attributed the shift of  $T_C$  to the internal stress existing in core-shell structure grains resulting from the thermal expansion mismatch between cores and shells. It is well accepted that the ferroelectric phase transition of BaTiO<sub>3</sub> is affected by stress, including external stress [1, 23] and internal stress [5, 11, 12, 22]. Curie temperature was shifted to lower temperatures by compaction external stress. Recently, the effect of internal stress on the Curie temperature of BaTiO<sub>3</sub> has been investigated. It was found by microindentation analysis that internal stress remaining in core-shell structure grains was much higher than that in chemical homogeneous grains, leading to a significant shift of Curie temperature [24]. Internal stress arose from the thermal expansion mismatch between grain cores and grain shells. Therefore, we examined the thermal expansion curves of pure BaTiO<sub>3</sub> and Nb-doped BaTiO<sub>3</sub> ceramics sintered at 1280 °C with various Nb<sub>2</sub>O<sub>5</sub> contents, as shown in Fig. 6. The thermal expansion coefficients calculated



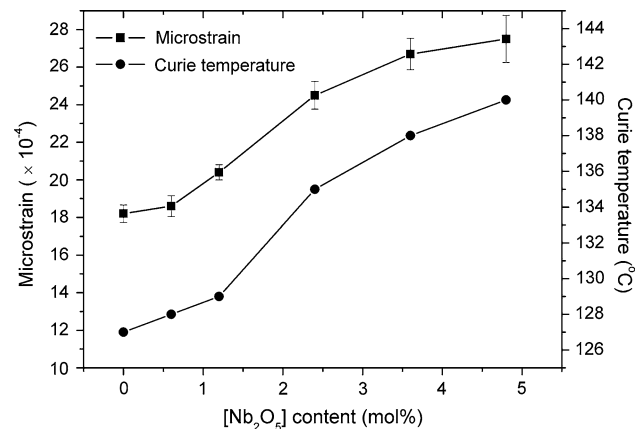
**Fig. 6** Thermal expansion curves of pure BaTiO<sub>3</sub> and Nb<sub>2</sub>O<sub>5</sub>-doped BaTiO<sub>3</sub> ceramics sintered at 1280 °C with various Nb<sub>2</sub>O<sub>5</sub> contents

**Table 1** Thermal expansion coefficient of pure BaTiO<sub>3</sub> and Nb<sub>2</sub>O<sub>5</sub>-doped BaTiO<sub>3</sub> ceramics sintered at 1280 °C

Nb <sub>2</sub> O <sub>5</sub> (mol%)	T <sub>0</sub> (°C)	ΔL <sub>0</sub> (μm)	T (°C)	ΔL <sub>t</sub> (μm)	L <sub>0</sub> (mm)	α <sub>L</sub> (ppm/°C)
0	50.0	7.2	99.0	23.0	51.98	6.772
0.3	51.2	2.5	100.1	19.1	53.58	6.906
0.6	51.1	6.3	100.2	23.4	54.18	6.998
1.2	50.8	3.7	100.2	22.8	53.44	7.805

from Fig. 6 are given in Table 1. It could be concluded that the thermal expansion coefficients of Nb-doped BaTiO<sub>3</sub> ceramics increased with an increase in Nb<sub>2</sub>O<sub>5</sub> contents. The ferroelectric core is pure barium titanate and the paraelectric shell was doped with additive Nb<sub>2</sub>O<sub>5</sub>. It was suggested that core and shell regions could not be totally separated in core-shell structure grains [3]. Therefore, it was reasonable to deduce that the thermal expansion coefficient of core-shell structure grains was a summation of those of cores and shells. Moreover, it was indicated from our results that the thermal expansion coefficient of shells was larger than that of pure barium titanate cores. According to high expansion of grain shell, pure barium titanate cores were exposed to compaction stress as cooling to room temperature. Moreover, the compaction stress placed on cores increased with increasing Nb<sub>2</sub>O<sub>5</sub> contents, due to an increase of the thermal expansion coefficient of grain shells.

On the other hand, according to Buessem et al. [25, 26], every grain in fine grained ceramics was subjected to a complex internal stress that depended on the orientation of all surrounding grains. From XRD analyses, a large amount of secondary phase grains existed with Nb<sub>2</sub>O<sub>5</sub> content greater than 1.2 mol%. Thus, it was considered that



**Fig. 7** Microstrain and Curie temperature as functions of Nb<sub>2</sub>O<sub>5</sub> contents for BaTiO<sub>3</sub> ceramics sintered at 1280 °C

Nb-doped BaTiO<sub>3</sub> grains were subjected to large internal stress resulting from the formation of a core-shell structure and secondary phase grains. In order to determine the microstrain of samples sintered at 1280 °C, the XRD patterns were subjected to detailed microstructure study by using MAUD. The results are reported in Fig. 7. The Curie temperature as a function of Nb<sub>2</sub>O<sub>5</sub> contents is also shown in Fig. 7. It could be seen that microstrain increased continuously with an increase in Nb<sub>2</sub>O<sub>5</sub> amount, and Curie temperature showed the same trend as microstrain. The Curie temperature moved from 127 to 140 °C in the case of 4.8 mol% Nb<sub>2</sub>O<sub>5</sub> addition. This was an indication that the Curie temperature of BaTiO<sub>3</sub> ceramics was dependent on microstrain to a certain extent.

## Conclusion

Homogeneous and fine grained microstructure was observed in Nb<sub>2</sub>O<sub>5</sub>-doped BaTiO<sub>3</sub> ceramics sintered at 1280 °C. For higher Nb<sub>2</sub>O<sub>5</sub> concentration ( $\geq 2.4$  mol%), large secondary phase grains occurred when sintered at 1320 °C. XRD analysis proved that secondary phase was detected when the Nb<sub>2</sub>O<sub>5</sub> amount exceeded 1.2 mol%, which was identified as Ba<sub>6</sub>Ti<sub>14</sub>Nb<sub>2</sub>O<sub>39</sub>, Ba<sub>3</sub>Ti<sub>5</sub>Nb<sub>3.2</sub>O<sub>21</sub>, and Ba<sub>6</sub>Ti<sub>17</sub>O<sub>40</sub>. It was found Nb<sub>2</sub>O<sub>5</sub> addition played an important role in the transition temperature of BaTiO<sub>3</sub> ceramics. As the amount of Nb<sub>2</sub>O<sub>5</sub> was increased, T<sub>C</sub> was shifted to higher temperatures. Contrarily, T<sub>1</sub> was independent of Nb<sub>2</sub>O<sub>5</sub> content up to 1.2 mol% and then decreased rapidly with further addition. The decrease of T<sub>1</sub> was assumed to be correlated with the secondary phase formation when Nb<sub>2</sub>O<sub>5</sub> concentration was greater than 1.2 mol%. On the other hand, the shift of T<sub>C</sub> to higher temperatures was attributed to large internal stress in Nb<sub>2</sub>O<sub>5</sub>-doped BaTiO<sub>3</sub> ceramics. The internal stress originates not only from the misfit between grain cores and

grain shells in core–shell structure, but also from the formation of the large secondary phase grains.

## References

1. Merz WJ (1953) *Phys Rev* 91:513
2. Lin JN, Wu TB (1990) *J Appl Phys* 68:985
3. Hennings D, Rosenstein G (1984) *J Am Ceram Soc* 67:249
4. Tang B, Zhang SR, Yuan Y. *J Mater Sci Mater Electron*. doi: [10.1007/s10854-007-9477-0](https://doi.org/10.1007/s10854-007-9477-0)
5. Jung YS, Na ES, Paik U (2002) *Mater Res Bull* 37:1633
6. Arlt G, Hennings D, de With G (1985) *J Appl Phys* 58:1619
7. Uchino K, Sadanaga E, Hirose T (1989) *J Am Ceram Soc* 72:1555
8. Begg BD, Vance ER, Nowotny J (1994) *J Am Ceram Soc* 77: 3186
9. Albertsen K, Hennings D, Steigelmann O (1998) *J Electroceram* 2(3):193
10. Lee S et al (2007) *J Appl Phys* 101:054119
11. Sato S, Fujikawa Y, Nomura T (2000) *Am Ceram Soc Bull* 79:155
12. Song YH, Hwang JH, Han YH (2005) *Jpn J Appl Phys* 44:1310
13. Kahn M (1971) *J Am Ceram Soc* 54:455
14. Cui B, Yu PF, Tian J (2007) *Mater Sci Eng A* 454–455:667
15. Sahu P, Pradhan SK, De M (2004) *J Alloys Compd* 377:103
16. Cont L et al (2002) *Ferroelectrics* 267:323
17. Iverson BJ, Jones JL, Bowman KJ (2006) *Phys B Condens Matter* 385–386:581
18. Brzozowski E, Castro MS, Foschini CR (2002) *Ceram Int* 28:773
19. Chiang SK, Lee WE, Readey DW (1987) *Am Ceram Soc Bull* 66:1230
20. Thomas NW (1990) *J Phys Chem Solids* 51:1419
21. Chazono H, Kishi H (2000) *J Am Ceram Soc* 83:101
22. Hwang JH, Choi SK, Han YH (2001) *Jpn J Appl Phys* 40:4952
23. Samara GA (1966) *Phys Rev* 151:378
24. Armstrong TR, Buchanan RC (1990) *J Am Ceram Soc* 73:1268
25. Buessem WR, Cross LE, Goswami AK (1966) *J Am Ceram Soc* 49:33
26. Buessem WR, Cross LE, Goswami AK (1966) *J Am Ceram Soc* 49:36